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# Consistency Test for the Determination of a Correct Phase Set of the Structure Factors

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When solving the phase problem by direct methods one usually has to decide which of the many suggested solutions is the correct one. The coefficients K and N defined in this paper can be calculated for each possible solution. A correct solution is expected to be that which gives the lowest values of the above coefficients.

#### Introduction

In recent years, considerable attention has been devoted to the problem of finding a correct phase set of the structure factors. Schenk (1972) criticized the shortcomings of the existing methods reflected mainly in the investigation of crystal structures in symmorphous space groups. Later, several procedures were published demonstrating how to avoid the above difficulties by using relationships making possible an indication of the negative sign of the structure invariant.

Hašek (1974) showed that these difficulties can also be removed in symmorphous space groups without applying the above special relationships. For the case of centrosymmetric structures he used a procedure in which the calculation of the coefficient K for the tested phase set of the structure factors is virtually reduced to the summation of triplets with the positive and with the negative sign respectively in the invidual groups according to their weights.

In this paper the procedure has been extended to include noncentrosymmetric space groups.

### **Theoretical relationships**

In recent years expressions have been derived for the mean values of cosine of the invariant  $\Phi_{111} = \varphi_H + \varphi_K + \varphi_{-H-K}$  (Karle & Gilardi, 1973) and  $\Phi_{211} = 2\varphi_H + \varphi_{K-H} + \varphi_{-H-K}$  (Giacovazzo, 1974). These expressions can be written as

$$\langle \cos \Phi \rangle_{\text{theor}} = \frac{I_1(w)}{I_0(w)} \simeq [(0.0106w - 0.1304)w + 0.5658]w.$$
(1)

The above approximation holds for the range  $0.0 \le w \le 6.0$ . For  $w \ge 6.0$ ,  $I_1(w)/I_0(w) \simeq 1$  (Main, 1971). The symbol  $I_i$  denotes a Bessel function of imaginary argument of the *i*th order. For the structure invariant  $\Phi_{111}$  the weight w is given by (Hauptman, 1972)  $w = 2 \cdot \sigma_3 \cdot \sigma_2^{-3/2} \cdot |E_H E_K E_{H-K}|$  or more exactly, (Karle & Gilardi, 1970),

$$w = 2\sigma_3 \cdot \sigma_2^{-3/2} \times \frac{|E_{II}E_KE_{H-K}|}{1+2|U_HU_KU_{H-K}| - |U_H|^2 - |U_K|^2 - |U_{H-K}|^2}, \quad (2)$$

where  $\sigma_i = \sum_{j=1}^{N} f_j^i$ ,  $f_j$  is the scattering factor of the *j*th atom,  $U_H$  is the unitary structure factor and N is the number of atoms in the unit cell. For the structure invariant  $\Phi_{211}$  the weight w is given by Giacovazzo (1974):

$$w = 2 \cdot \sigma_3^2 \cdot \sigma_3^2 (2 \cdot |E_K|^2 - 1) |E_H^2 \cdot E_{K-H} E_{-H-K}| .$$
(3)

For the variances of the cosine of the invariant  $\Phi_{111}$ and  $\Phi_{211}$  relationships have been derived (Hauptman, 1972; Karle & Gilardi, 1973; Giacovazzo, 1974) which can be written as

var (cos 
$$\Phi_{\text{theor}}$$
) = 1 -  $\frac{I_1(w)}{w \cdot I_0(w)} - \left(\frac{I_1(w)}{I_0(w)}\right)^2$  (4)

where the weight w is given by (2) for the invariant  $\Phi_{111}$ and by (3) for the invariant  $\Phi_{211}$ . It can be expected for structures with a great number of atoms in the unit cell that the mean values  $\langle \cos \Phi \rangle_{calc}$  and the variances var  $(\cos \Phi)_{calc}$  calculated for a correct phase set of the structure factors will be very well approximated by the corresponding values calculated by using (1) and (4). In the case of an incorrect phase set, when on the respective Fourier map there may appear a region with negative values of the 'electron density' or exceedingly strong false peaks, the correspondence between the  $\langle \cos \Phi \rangle_{calc}$  and var  $(\cos \Phi)_{calc}$  values on the one hand and the respective  $\langle \cos \Phi \rangle_{theor}$  and var  $(\cos \Phi)_{theor}$  values on the other will be considerably poorer.

This principle is the basis of the following consistency test.

## **Consistency test**

The structure invariants  $\Phi$  were divided into p groups so that each group contained a sufficiently high number of invariants of approximately the same weight. For each group the mean weight  $\bar{w}^{(l)}$  was calculated, and the corresponding mean values of  $\langle \cos \Phi \rangle_{\text{theor}}^{(l)}$  and the variances var  $(\cos \Phi)_{\text{theor}}^{(l)}$  were determined by using (1) and (4) respectively.

The values of the structure invariants  $\Phi_{cale}$  were calculated for each phase set of the structure factors tested. Further, for each group of invariants the average value  $\langle \cos \Phi \rangle_{cale}^{(i)}$ , an estimate of the variance of

$$\operatorname{var}(\cos \Phi)_{\operatorname{calc}}^{(i)} = \langle (\cos \Phi)^2 \rangle_{\operatorname{calc}}^{(i)} - (\langle \cos \Phi \rangle_{\operatorname{calc}}^{(i)})^2 \rangle_{\operatorname{calc}}^{(i)} = \langle (\cos \Phi)^2 \rangle_{\operatorname{c$$

and an estimate of the variance of the mean value of the cosine of the invariant var  $(\langle \cos \Phi \rangle)_{calc}^{(i)}$  were also calculated.

It may be expected that the correct phase set of the structure factors will be found among the sets giving the best fit between the values of  $\langle \cos \Phi \rangle_{\text{theor}}^{(l)}$  and  $\langle \cos \Phi \rangle_{\text{cale}}^{(l)}$ . If we assume that  $\langle \cos \Phi \rangle_{\text{cale}}^{(l)}$  values are statistically independent, a suitable measure of the agreement between these values is the weighted sum of their deviations squared

$$K = \frac{\sum_{i=1}^{p} r_i (\langle \cos \Phi \rangle_{\text{calc}}^{(i)} - \langle \cos \Phi \rangle_{\text{theor}}^{(i)})^2}{\sum_{i=1}^{p} r_i}, \quad (5)$$

where

$$r_l = [\operatorname{var}(\langle \cos \Phi \rangle)_{\operatorname{calc}}^{(l)}]^{-1}.$$
(6)

At the same time, for the correct phase set of structure factors the coefficient

$$N = \sum_{l=1}^{p} [\text{var}(\cos \Phi)_{\text{calc}}^{(l)} - \text{var}(\cos \Phi)_{\text{theor}}^{(l)}]^2, \quad (7)$$

also must assume one of the least values. For the invariant  $\Phi_{111}$  the coefficients K and N will be denoted by the symbols  $K_{111}$  and  $N_{111}$ , while for the invariant  $\Phi_{211}$  the respective notation will be  $K_{211}$  and  $N_{211}$ .

Thus, the most probable phase system of structure factors can be seen in a system giving the lowest values of the coefficients  $K_{111}$  and  $N_{111}$  (or  $K_{211}$  and  $N_{211}$ ).

Similarly to the case of the coefficient K (Hašek, 1974), for centrosymmetric space groups, a connexion between the coefficient  $K_{111}$  and the absolute figure of merit (Germain, Main & Woolfson, 1971) can also be pointed out. If only one weight interval is taken in the calculation of the coefficient  $K_{111}$ , *i.e.* (5) is reduced to

$$K = \langle \cos \phi \rangle_{\text{calc}} - \langle \cos \phi \rangle_{\text{theor}}$$
,

then the coefficient  $M = |1 - K_{111}|$  will have properties similar to those of the absolute figure of merit.

The reliability of the consistency test depends on the

conditions of the validity of (1) and (4). In the case of great overlapping of the Patterson peaks, pseudosymmetry, or if the dominant effect on diffracted intensities has a low number of heavy atoms, it may occur that the correct phase set of the structure factors will not correspond to the lowest values of the coefficients K and N, but only to some of the following values. Unreliably determined values of normalized structure factors can have a similar effect.

### Conclusion

There are many routes to the determination of a correct phase set of the structure factors. The advantage of the procedure just outlined compared to the majority of the other methods is seen in its simplicity: calculations using relationships (1), (2), and (5) consist only in a correct classification and averaging of values already calculated by the common direct methods.

Since the procedure used here for the calculation of the coefficients  $K_{111}$  and  $N_{111}$  – simple as it is – suppresses the frequently criticized shortcomings of the figures of merit of symmorphous space groups, it seems to be a suitable supplement to the majority of the existing direct methods employing triplet relationships for the investigation of crystal structures.

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